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Cosmic Rays

Alan Watson University of Leeds

a.a.watson@leeds.ac.uk

OUTLINE

Lecture 1

Some History

Properties of Low Energy Cosmic Rays

Tutorial on physics of air showers

Cosmic rays at $\approx 10^{15} \text{ eV}$ – the 'knee'

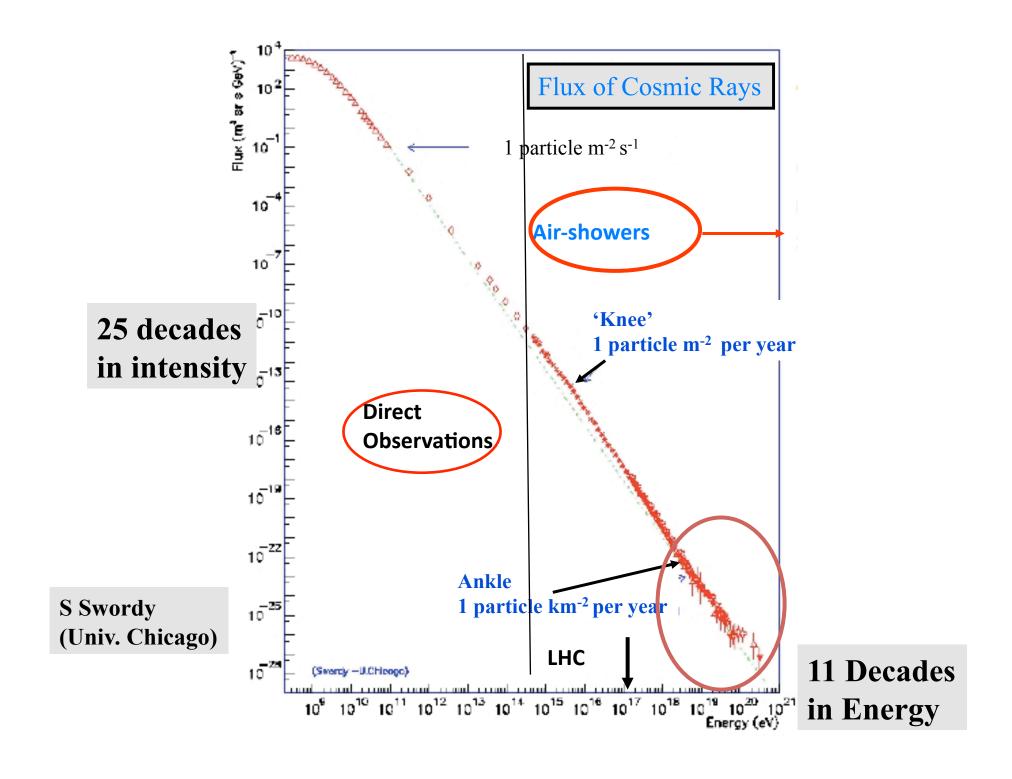
Lecture 2

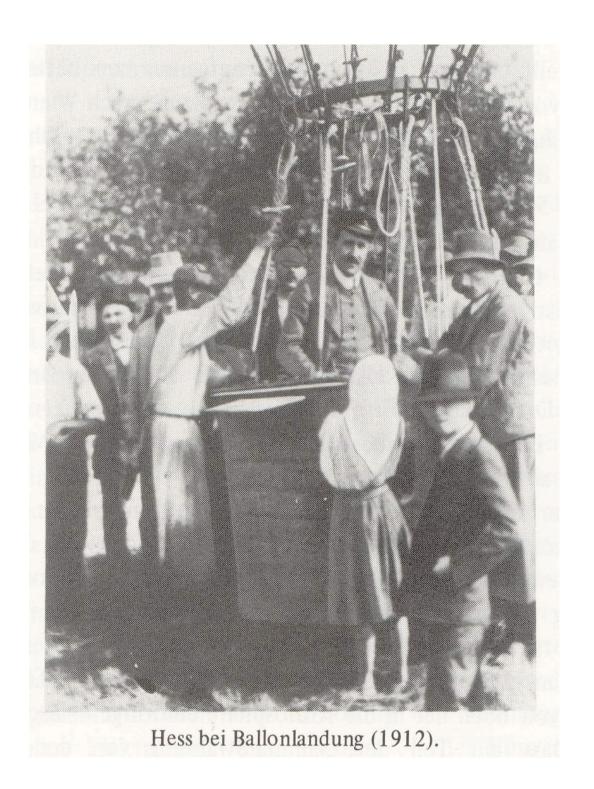
The Highest Energy Cosmic Rays

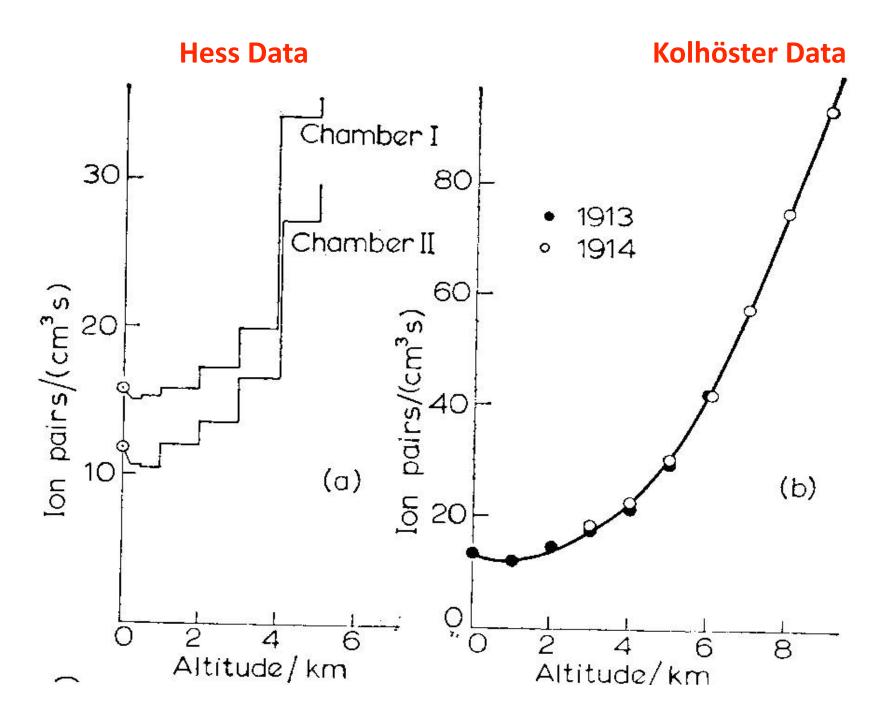
Why study them?

Pierre Auger Observatory

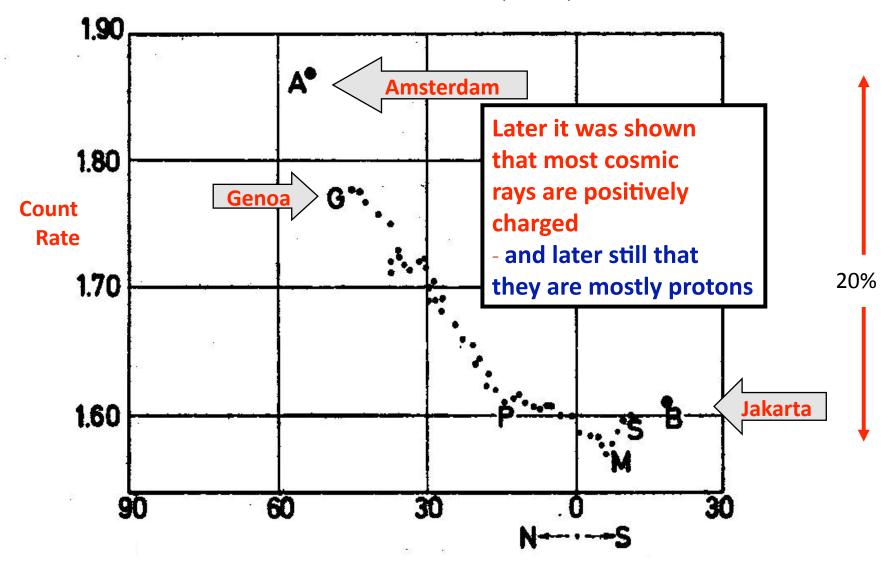
Results



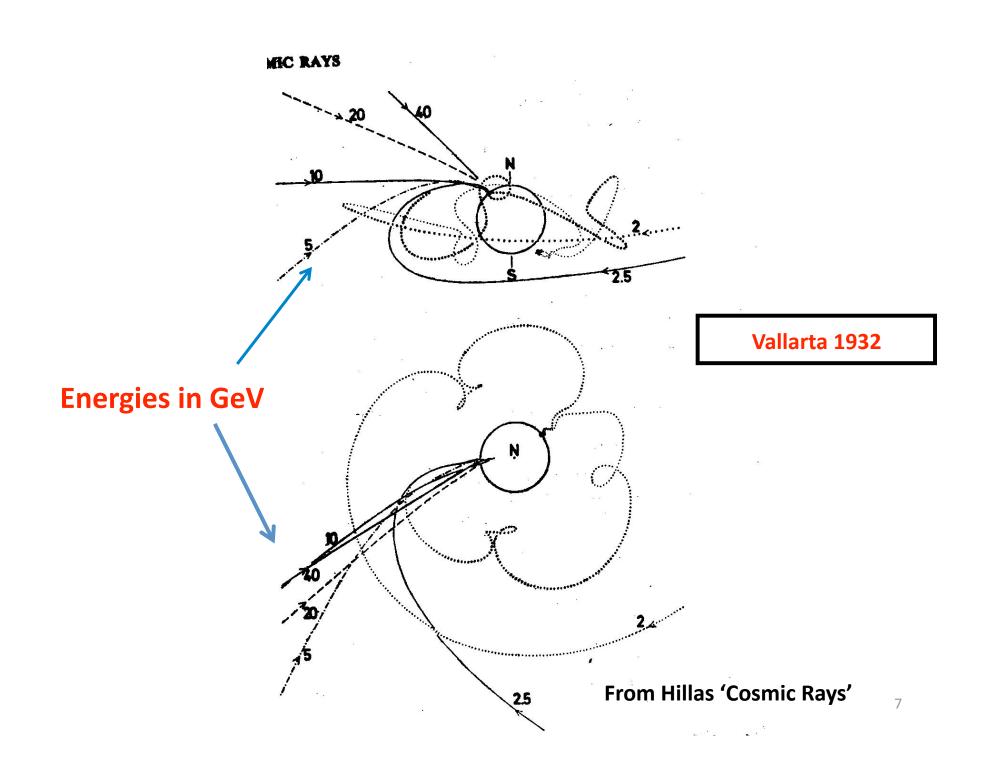




Clay's Results, taken by Berlage (~1926)



Counting Rate at sea-level as a function of the position of the ship with respect to the earth's magnetic field which is nearly horizontal at the earth's equator

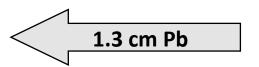


Some properties of low energy cosmic rays

- Energy density of cosmic rays ≈ 1 eV cm⁻³
- Similar to energy density in 2.7 K radiation, starlight, turbulent gas motions and magnetic fields only the last two are significant equalities
- 1 eV cm⁻³ through out Universe faces us with a daunting problem
- For low energy cosmic rays, our galaxy is a place where sources are likely to be located: for the highest energy cosmic rays???
- Can be observed directly up to ≈ 100 TeV from balloons (e.g. ATIC) and from space (e.g. AMS watch out for results from this)
- Detectors are used to measure –dE/dx and the charge, with the energy being measured in a variety of ways, often with a calorimeter
- Electrons≈ 10⁻² and gamma rays ≈ 10⁻⁴

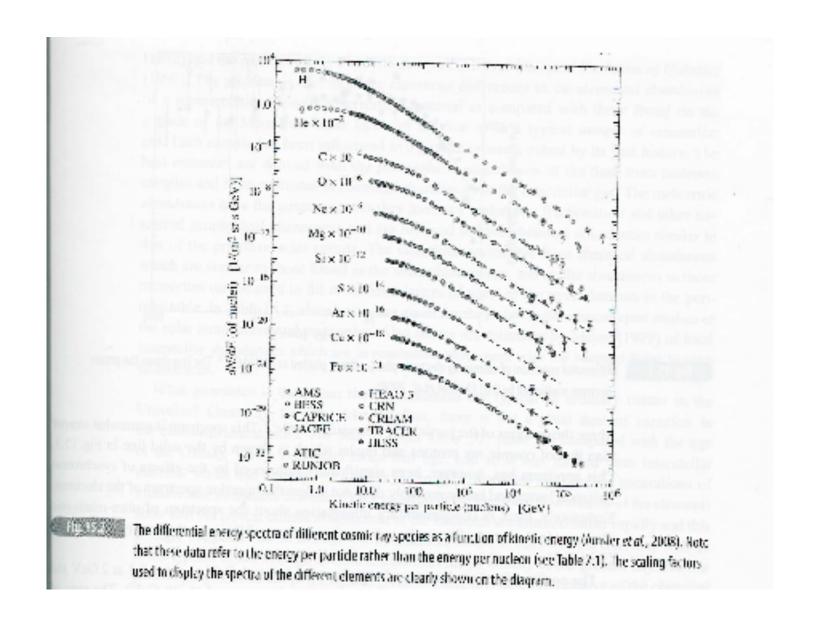
(Im)Practical example of how it is done

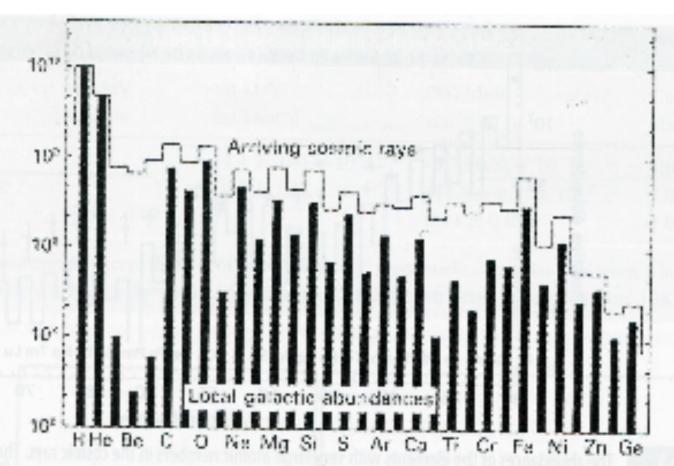




- Incoming particle is highly likely to be a proton
- Level of ionisation excludes heavier nucleus (dE/dx) ~Z²
- Traversal of the particle through 6 Pb-plates (about 88.5 g cm⁻² or 13.9 rad. lengths)

without interaction strongly excludes an electron.





The cosmic abundances of the elements in the cosmic rays (solid line) compared with the Solar System abundances isolid histogram). The data have been normalised to a relative abundance of hydrogen of 10% (Lund, 1984).

Longair p 495

The comparison of the distribution of the elements shows some striking features

- Presence of Li, Be and Boron in relatively large quantities
- Presence of elements just lighter than Fe

Ratio of ³He to ⁴He (needs sensitivity to isotopes)

Spallation:

Interaction of parent nuclei with hydrogen of ISM

See analysis in Longair pp 507 to 517

From the observed abundances, using a set of transfer equations, Longair pp 507 et seq

One finds that typically these cosmic rays have traversed ≈ 5 g cm⁻²

Assuming that there is 1 atom of H per cm³

This sets the lifetime of the cosmic rays at 3×10^6 to 3×10^7 years

This age can be confirmed by using various radioactive nuclei, such as 10 Be as a radioactive clock ($\tau = 1.51 \times 10^6$ years)

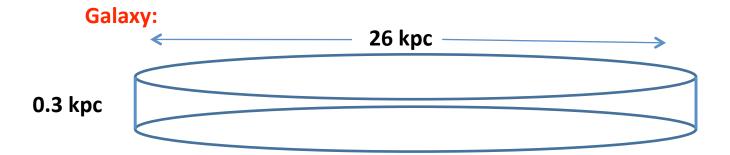
Study of the transport equations is also important for appreciating manner in which electrons propagate: study them

- Direct measurements are possible up to about 100 TeV
 e.g. ATIC experiment
- In terms of energy per nucleus

p: He: 2 x CNO: 2 x Ne-Si : 2 x Z>17 : 4 x Fe

• Lifetime, energy density and storage volume give us an estimate of the power of the cosmic ray sources

We need to assume the storage volume and the indications from the lifetime points to the galactic plane

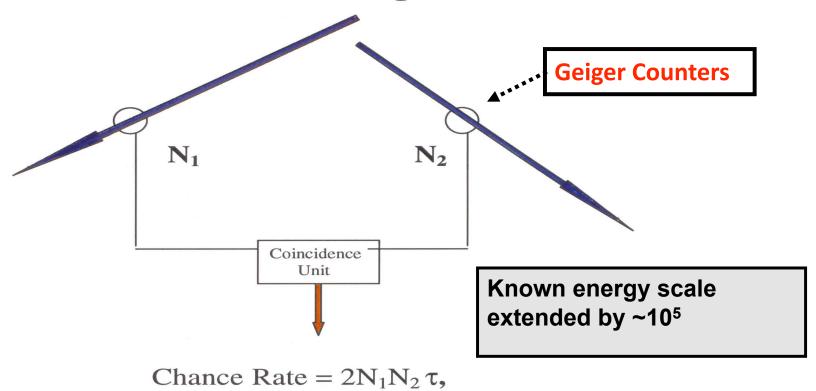


- Total energy in this volume is about 5 x 10⁵⁴ ergs
- Ginsburg in 1960s argued that Supernovae (SN)
 could power the cosmic rays
- 1 per 30 years in the galaxy (to within about x2)
- So about 2.5 x 10⁴⁹ ergs in cosmic rays from each SN
- This is about the same as the energy in visible light

and about 2% of kinetic energy output of SN

(You will hear discussion about evidence of whether or not SN ARE the source of cosmic rays from Jim Hinton)

Discovery of Extensive Air Showers: Pierre Auger (1938)



Resolving time = 10 millionth of a second (10^{-5} s)

Observed Rate was found to be much higher than the Calculated Chance Rate

 even when the counters were as far as 300 m apart (at Jungfrauhoch)



~ 1 TeV

Fretter: Echo Lake, 1949

Reasonable to have point of interaction in 7th plate (this depends on p- Pb cross-section)

$$p_{CR} + p \rightarrow p + p (or n) + N(\pi^{+} + \pi^{-} + \pi^{0})$$

Also K, Λ , η , Ω , Σ are undoubtedly created

What is the energy of the particle?

- a MUCH harder question to answer in larger showers

How can we understand shower development?

Electromagnetic part of the cascade (i)

Key processes are bremsstrahlung and pair production: cross-sections calculable from QED

$$\lambda_{\mbox{\tiny pair}}$$
 = 1/n $\sigma_{\mbox{\tiny pair}}$, where

$$\sigma_{\text{pair}} \cong ((Z^2 r_e^2)/137)(28/9)\ln(183/Z^{1/3}) \text{ cm}^2$$

or
$$\sigma_{pair} \cong 5.7 \, r_e^2 = 6 \times 10^{-26} \, cm^2$$
, for air

$$\sigma_{\text{\tiny pair}} = (7/9)\sigma_{\text{\tiny brem}}$$

Electromagnetic part of the cascade (ii)

- gamma ray disappears in pair production
- opening angle: $\theta \approx mc^2/hv$

- bremsstrahlung energy spectrum is 'flat'
 - electron can lose all of its energy at once
- opening angle ≈mc²/E, where E is energy of radiated photon

The Radiation Length

 $E = E_o \exp(-x/X_o)$, where X_o is the radiation length

if,
$$-dE/dx)_{brem} = -dE/dx)_{ionisation}$$

electron is said to be at the critical energy, E = ε_c

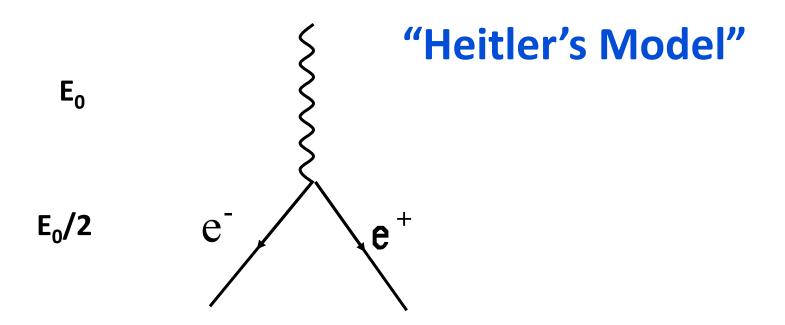
(Good source for derivations of relations discussed in last few slides: Longair 'High Energy Astrophysics', Second Edition, Volume 1) The following discussion and slides are due to Jim Matthews (LSU). The treatment is an approximation intended to exhibit some of the physics driving the main features of air showers plainly and simply

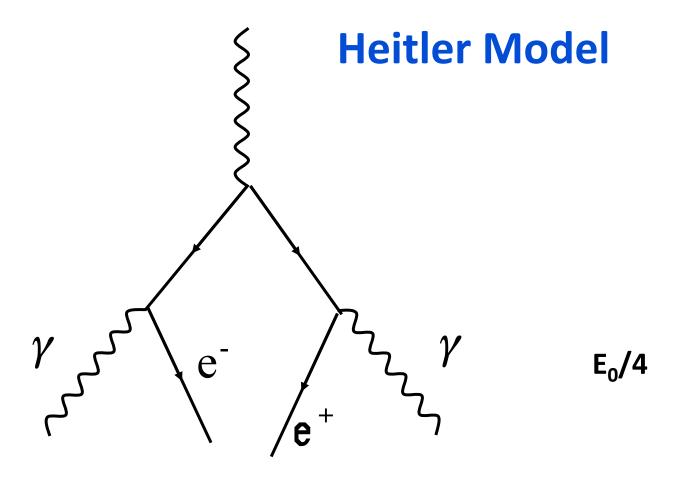
It does not replace full simulations

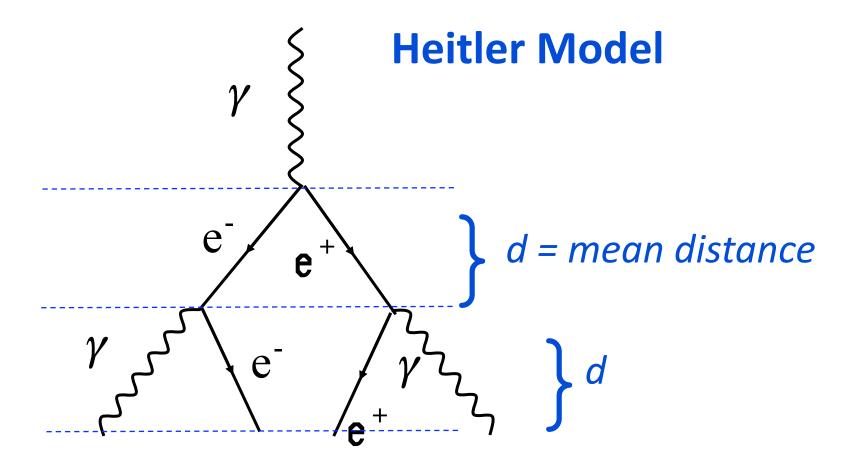
It is a useful pedagogic tool from which we can learn a lot

J Matthews: Astropart. Phys. 22 (2005) 387.

 $\begin{cases} & \text{The Heitler Model} \\ \gamma & \end{cases}$







$$d = \lambda_r \ln 2$$

 $\lambda_r = 37 \text{ g cm}^{-2} \text{ in air}$ $(radiation length)_{26}$

After "n" splits, there are "N" particles (e+, e-, and photons):

$$N = 2^n = e^{x/\lambda_r}$$

After "n" splits, there are "N" particles (e⁺, e⁻, and photons):

$$N = 2^n = e^{x/\lambda_r}$$

- Energy is evenly split between two secondaries.
- Cascade stops after " n_c " splits when individual energies are too low: *critical energy* ξ_c .

 $(\xi_c \text{ is when collision losses} >$ radiative losses: 85 MeV in air)

$$N_{max} = 2^{n_c}$$

$$E_{\circ} = \xi_c^e N_{max}$$

$$n_c = \ln[E_{\circ}/\xi_c^e]/\ln 2$$

Things the Heitler Model does well:

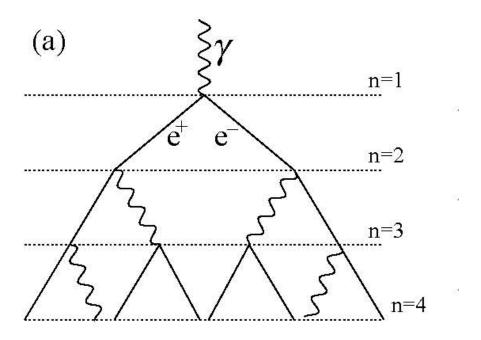
 $N_{max} \sim E_o$ - but not constant of proportionality

$$X_{\text{max}} \sim \log E_{\text{o}}$$

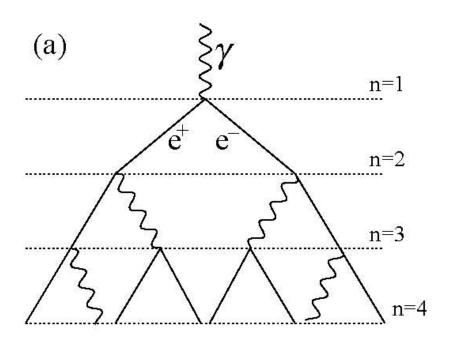
$$\Lambda \equiv rac{{
m d}\, X_{max}}{{
m d}\log_{10}E_{\circ}}~$$
 = 2.3 $\lambda_{
m r}$ = (85 g cm $^{ ext{-2}}$)/decade

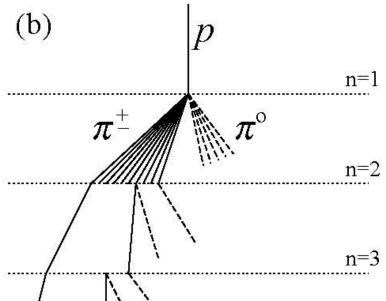
Things it does not do:

- relative numbers of photons/electrons
- attenuation, especially after maximum



$$E_{\circ} = \xi_c^e N_{max}$$





Extension to hadronic cascade is Jim Matthews's original contribution

$$E_{\circ} = \xi_c^e N_{max}$$

$$N_{ch} = N_{\pi \pm} = 10$$
 , $N_{\pi o} = 5$
$$N_{\mu} = N_{\pi \pm}$$

$$E_{\circ} = \xi_c^e N_{max} + \xi_c^{\pi} N_{\mu}$$

For pions, the distance between events is defined by the *interaction length*

$$d = \lambda_1 \ln 2$$

$$\lambda_1 = 120 \text{ g cm}^{-2}$$

(c.f.
$$\lambda_p = 80 \text{ g cm}^{-2}$$
)

The hadronic *critical energy* is reached when the distance to the next interaction exceeds the (dilated) lifetime

$$\xi_c$$
 = 20 GeV

After n generations, there are N_{π} charged pions:

$$N_{\pi} = (N_{ch})^n$$

Total energy carried by all these pions:

$$(2/3)^n E_o$$

e.g. ~ 10% after 5 generations

So each pion has:

$$E_{\pi} = \frac{E_{\circ}}{(\frac{3}{2}N_{ch})^n}$$

When pions drop below their critical energy:

$$E_{\pi} = \xi_{c} = 20 \text{ GeV}$$

all π[±] decay to muons

$$\ln N_{\mu} = \ln N_{\pi} = n_c \ln N_{ch} = \beta \ln [E_{\circ}/\xi_c^{\pi}]$$

$$\beta = \frac{\ln[N_{ch}]}{\ln[\frac{3}{2}N_{ch}]} = 0.85$$

$$N_{\mu} = \left(\frac{E_{\circ}}{\xi_c^{\pi}}\right)^{\beta} \approx 10^4 \left(\frac{E_{\circ}}{1 \text{ PeV}}\right)^{0.85}$$

(Full simulations give $\beta = 0.85 - 0.92$)

(n.b.: logarithmic dependence on N_{ch})

The growth of N_{μ} with E_o is less-than-linear (θ < 1).

Lower energy showers are more "efficient" in muon production

This is why Fe primaries make more muons than protons do (superposition model: 56 showers each with $E = E_p/56$)

8 depends on the (logarithmic) ratio of charged to neutral pions

The primary energy of the shower is divided into EM and hadronic channels:

$$E_{\circ} = \xi_c^e N_{max} + \xi_c^{\pi} N_{\mu}$$

Use observed $N_e = N_{max}/g$, $g \approx 10$:

$$E_{\circ} = g\xi_c^e(N_e + \frac{\xi_c^{\pi}}{g\xi_c^e}N_{\mu}) \approx 0.85 \text{ GeV}(N_e + 24N_{\mu})$$

A great deal can be learned by measuring N_e and N_μ in the same events

Depth of shower-maximum must be treated a little more carefully because it strongly depends on the *first* interaction.

- 1. Do an EM shower with $(1/3 E_o)/N_{ch}$
- 2. Use increasing multiplicity $N_{ch} \sim E_o^{1/5}$
- 3. Use energy dependent p-air λ_i

$$X_{max}^{p} = X_{\circ} + \lambda_{r} \ln[E_{\circ}/(3N_{ch}\xi_{c}^{e})]$$

= $(470 + 58 \log_{10}[E_{\circ}/1 \text{ PeV}]) \text{ g cm}^{-2}$

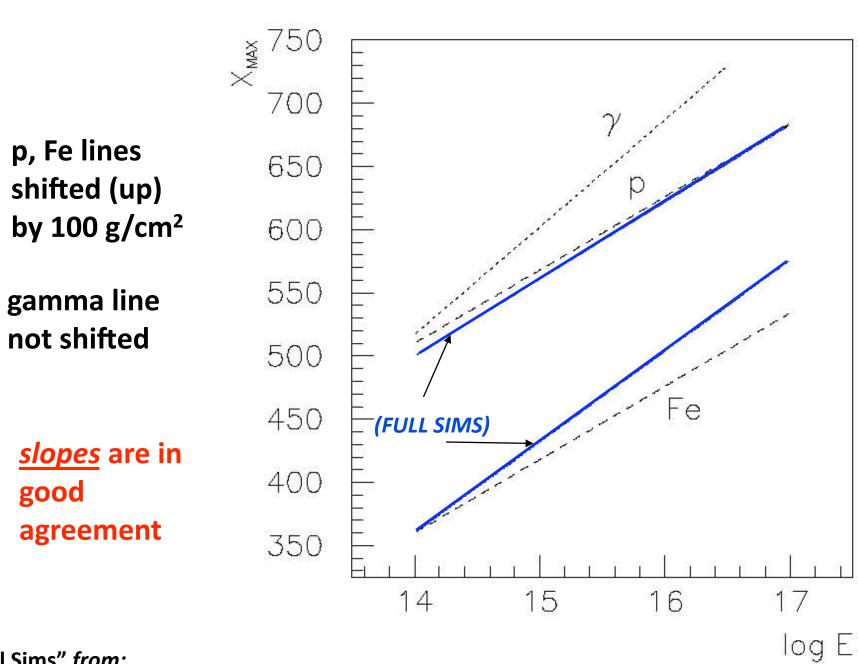
Not deep enough by $\approx 100 \text{ g cm}^{-2}$

Express in terms of EM-shower X_{max} :

$$X_{max}^p = X_{max}^\gamma + X_\circ - \lambda_r \ln[3N_{ch}]$$

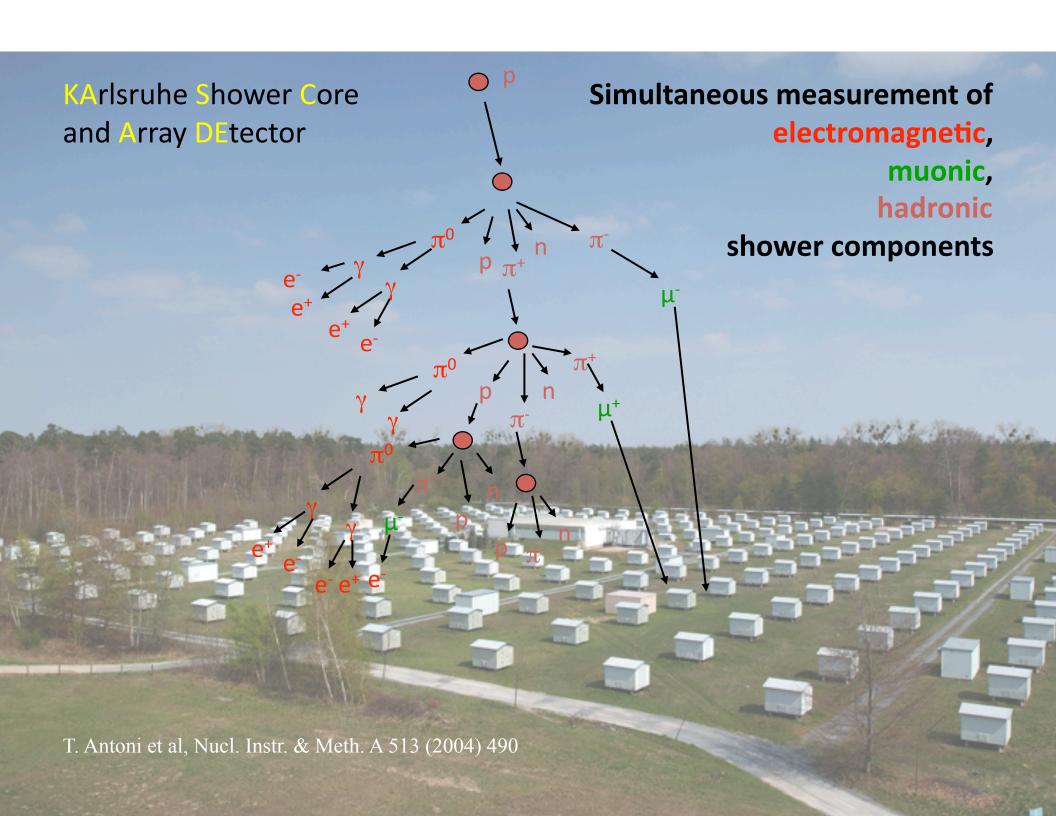
Elongation rate is in very good agreement with detailed simulations:

$$\Lambda^p = \Lambda^{\gamma} + \frac{\mathrm{d}}{\mathrm{d}\log_{10} E_{\circ}} \{X_{\circ} - \lambda_r \ln[3N_{ch}]\} = 58 \text{ g cm}^{-2} \text{ per decade}$$

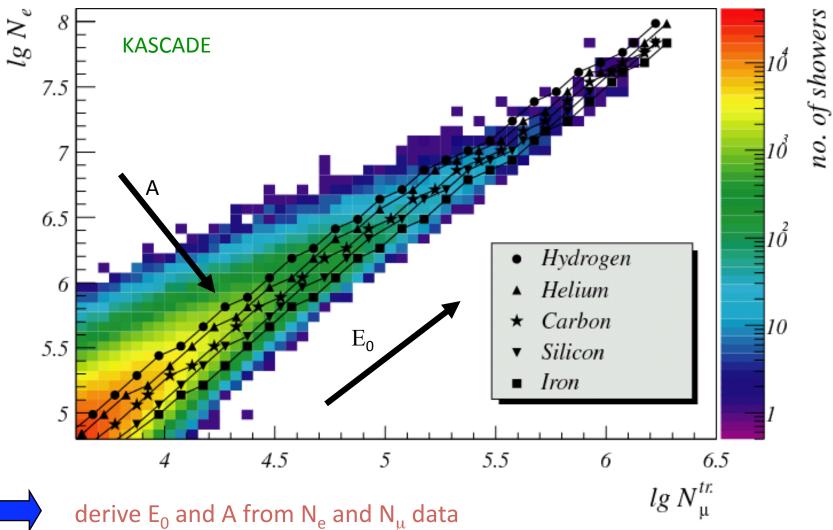


"Full Sims" from:

40



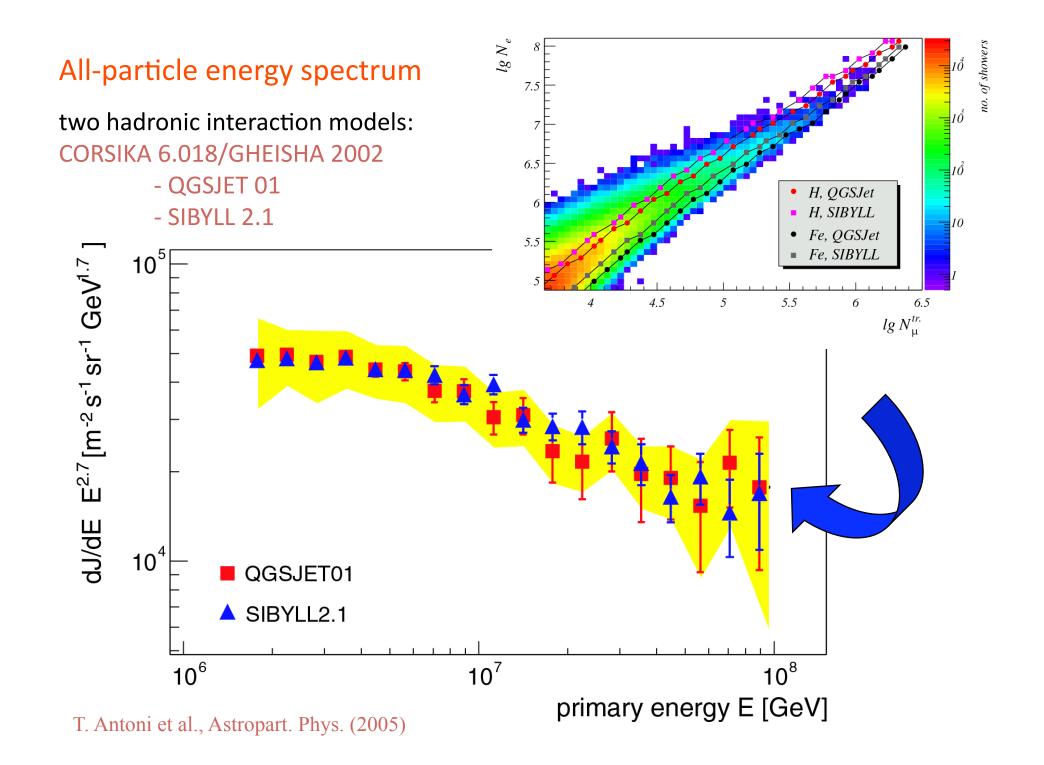
Two dimensional shower size spectrum $\lg N_e$ vs. $\lg N_{\mu}$

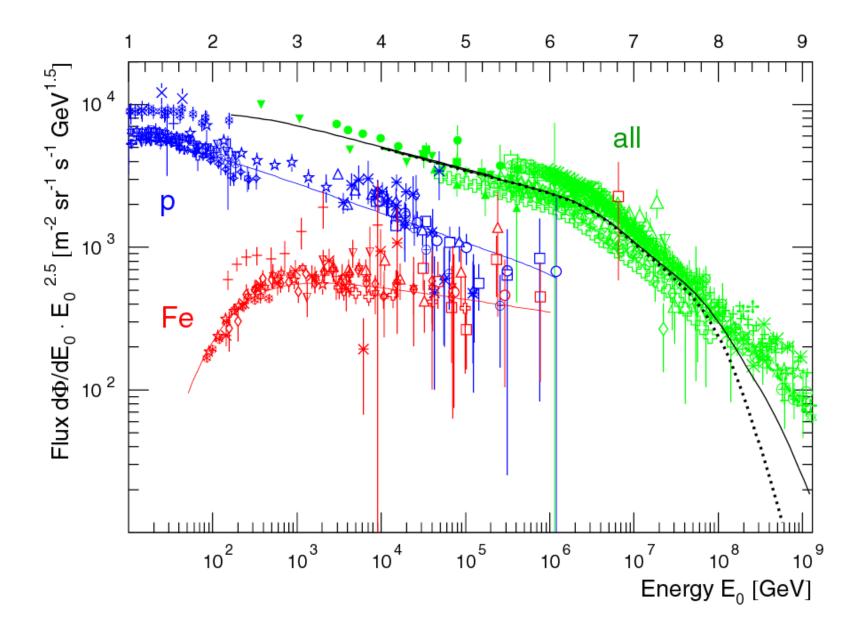




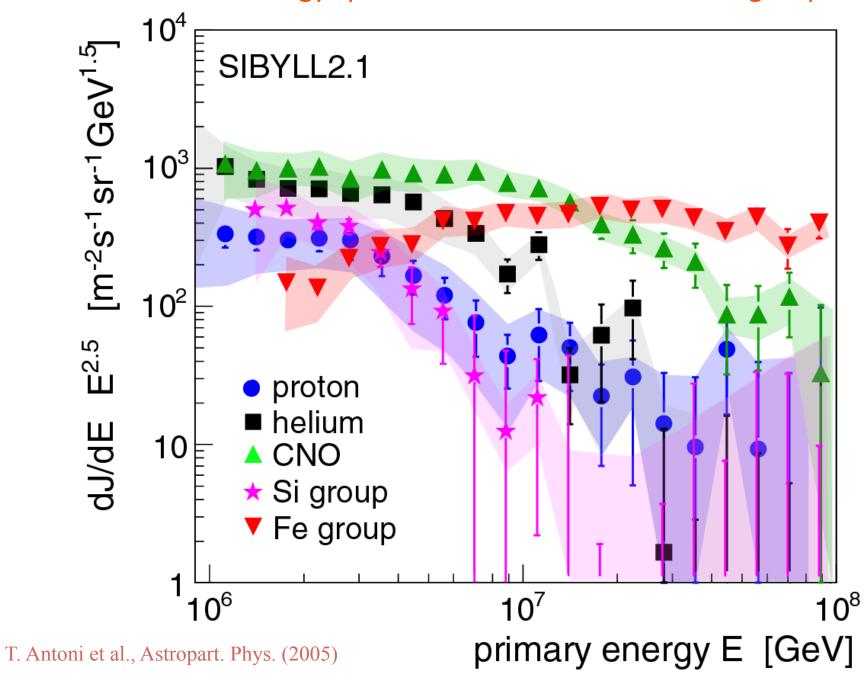
Fredholm integral equations of 1st kind:

$$g_i(\lg N_e, \lg N_\mu) = \int_0^\infty t_i(\lg N_e, \lg N_\mu \mid E) p_i(E) dE$$





KASCADE: Energy spectra for individual elemental groups



For many years there was a huge debate as to whether or not the bend in the spectrum, the knee, was due to particle physics (bend at energies increasing by A) or bend was at energy increasing by Z

The KASCADE results strongly suggest that the spectrum is changing because of a rigidity effect – change depends on Z

But not yet clear that this is a source or a leakage effect

In my view different types of sources can't be excluded and that's the end of part I



Giving a golf ball the energy of a high-energy cosmic ray: South Pole, Jan 1988

Why study Ultra-High Energy Cosmic Rays?

- no idea of their origin
- how to accelerate to 10²⁰eV?
- steepening of spectrum at highest energies?

 PROPAGATION EFFECTS or SOURCES?

Difficulties:

Above 10¹⁹ eV the rate is ~ 1 km⁻² per year

- energies are hard to measure
- mass spectrum is unknown
- anisotropies are hard establish

Spectrum shape:

$$E = 2 \Gamma ε2.7 K$$
 (for head-on collision)

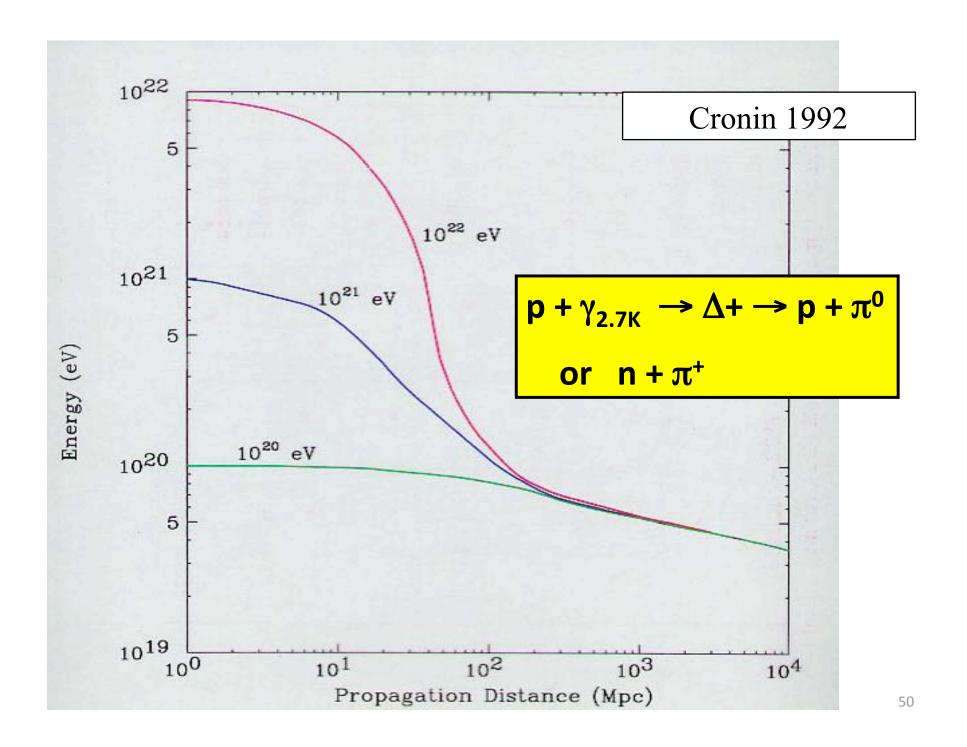
Steepening above 4 x 10¹⁹ eV? (GZK-effect: 1966)

$$\gamma_{2.7~K} + p \rightarrow \Delta^{+} \rightarrow n + \pi^{+} \text{ or } p + \pi^{-} \text{ (CMB well-known)}$$

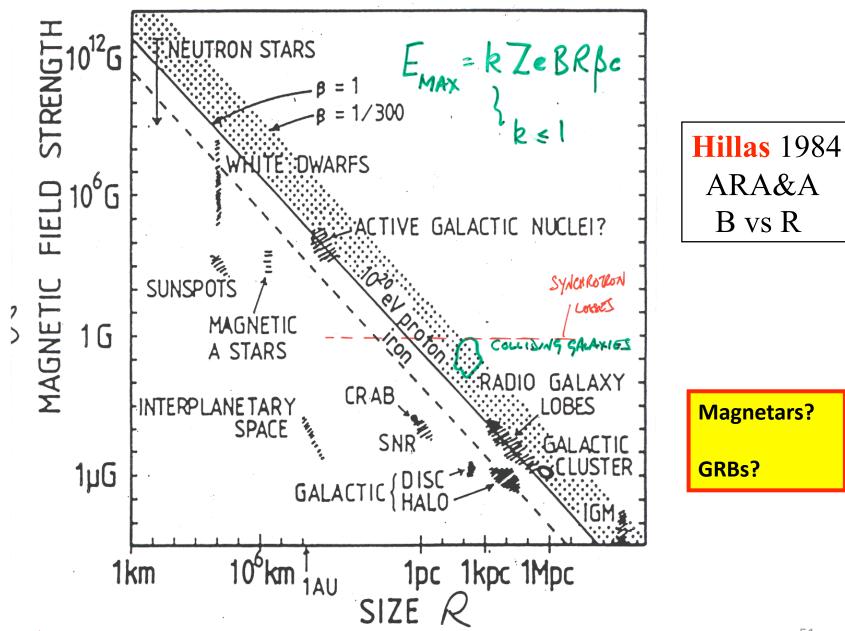
$$\gamma_{IR} + A \rightarrow (A-1) + n$$
 (IR background poorly known)

Also
$$\gamma_{2.7K} + Z \rightarrow Z + e^+ + e^-$$
 (pair production)

$$\gamma + \gamma_{\text{radio}} \rightarrow e + e$$
 (but 100 MHz background at UHECR unknown)



THE HILLAS PLOT (ANN ROV AS. Ap 19824)



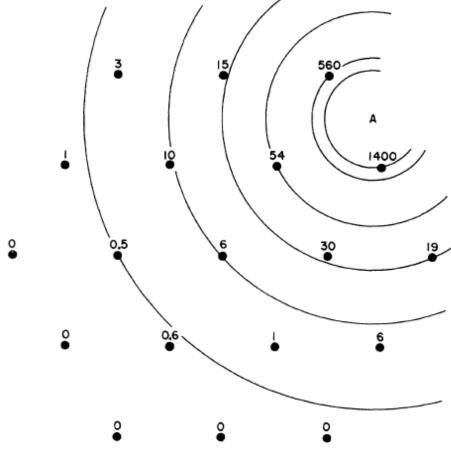
Particles in region of predicted GZK-steepening could tell us about sources within 100 – 200 Mpc - depending on the energy.

IF particles are protons, the deflections are expected to be small enough above $\sim 5 \times 10^{19}$ eV ($\sim 2^{\circ}$) that point sources might be seen – provided there are not too many.

So, measure:

- energy spectrum to look for GZK-prediction
- arrival direction distribution explore
- mass composition for interpretation

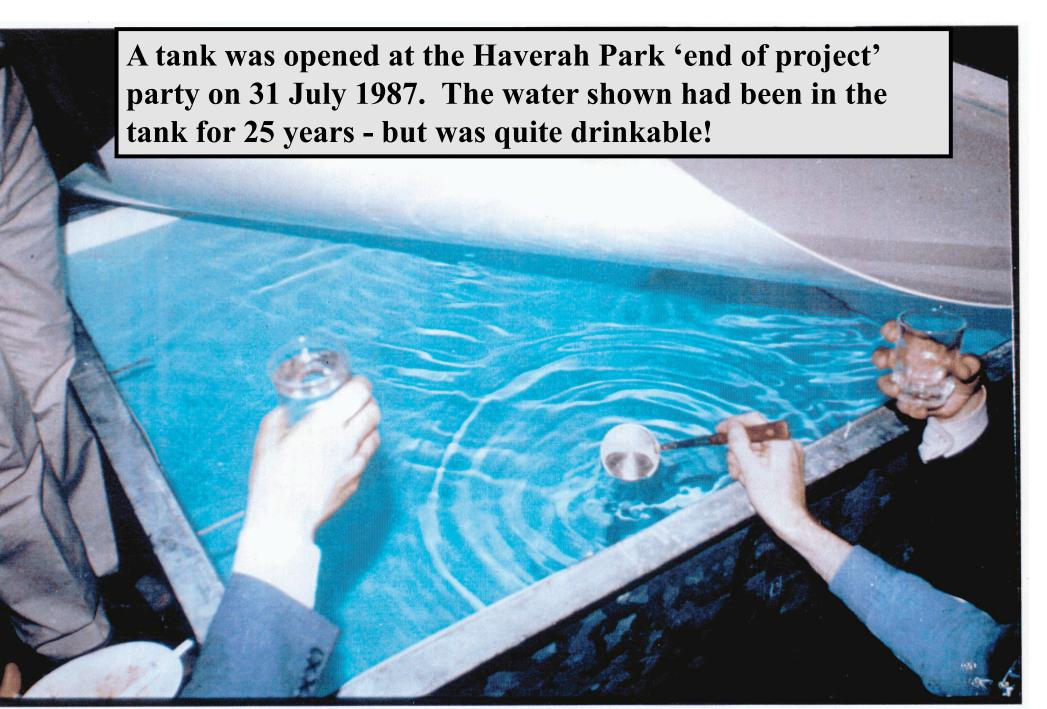




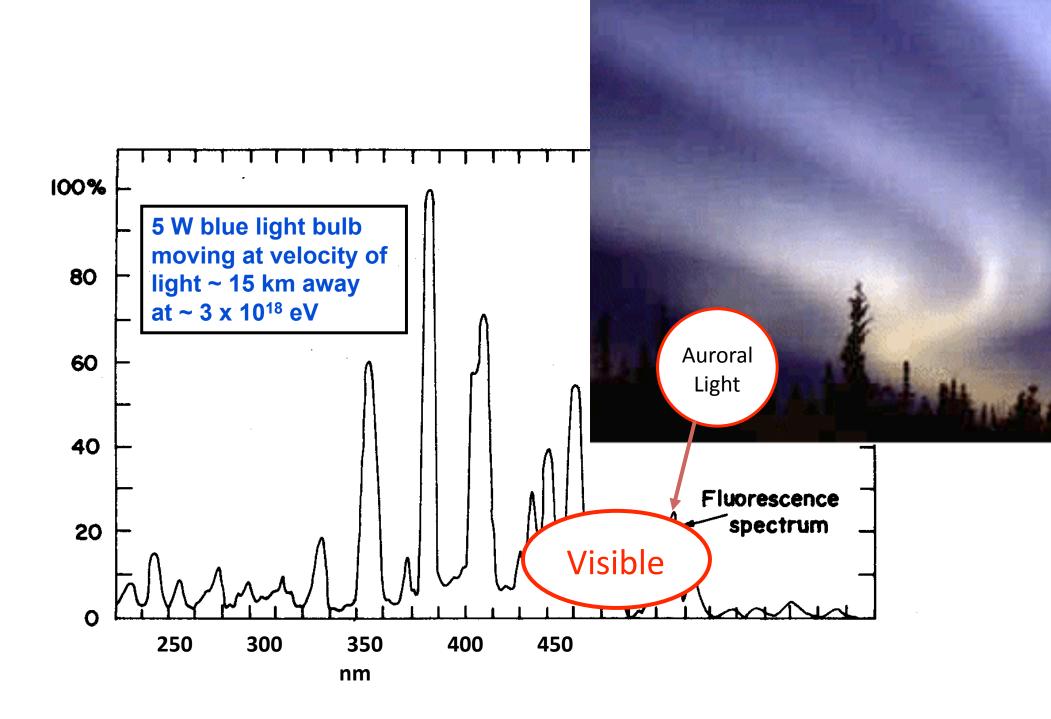
John Linsley (1927 – 2002)

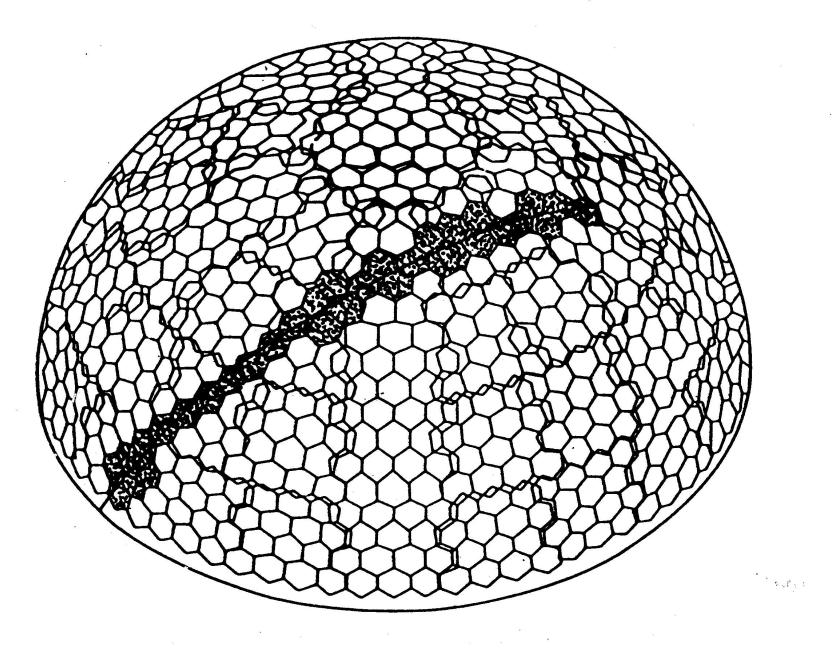
Pioneer of Large Shower Arrays

Volcano Ranch: 10²⁰ eV





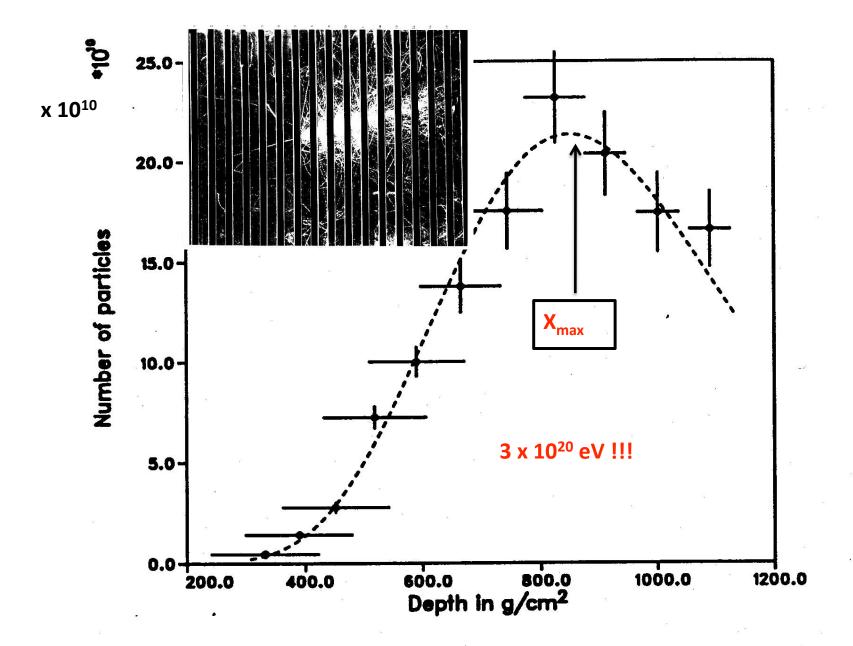




Idea of Fly's Eye Detector (University of Utah): 880 photomultipliers

HiRes: detector of fluorescence light





The Design of the Pierre Auger Observatory marries the two techniques just described in

the 'HYBRID' technique

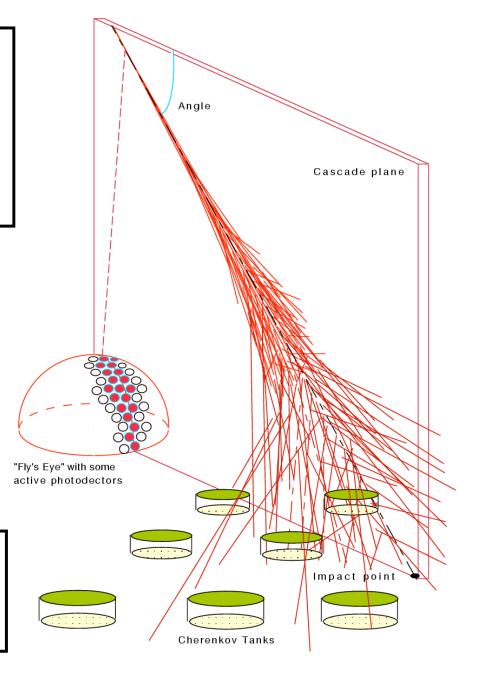
Fluorescence

AND

Array of water-

Cherenkov detectors





The Pierre Auger Collaboration

Czech Republic Argentina

France Australia

Germany Brasil

Italy Bolivia*

Netherlands Mexico

Poland USA

Portugal Vietnam*

Slovenia

Spain

United Kingdom

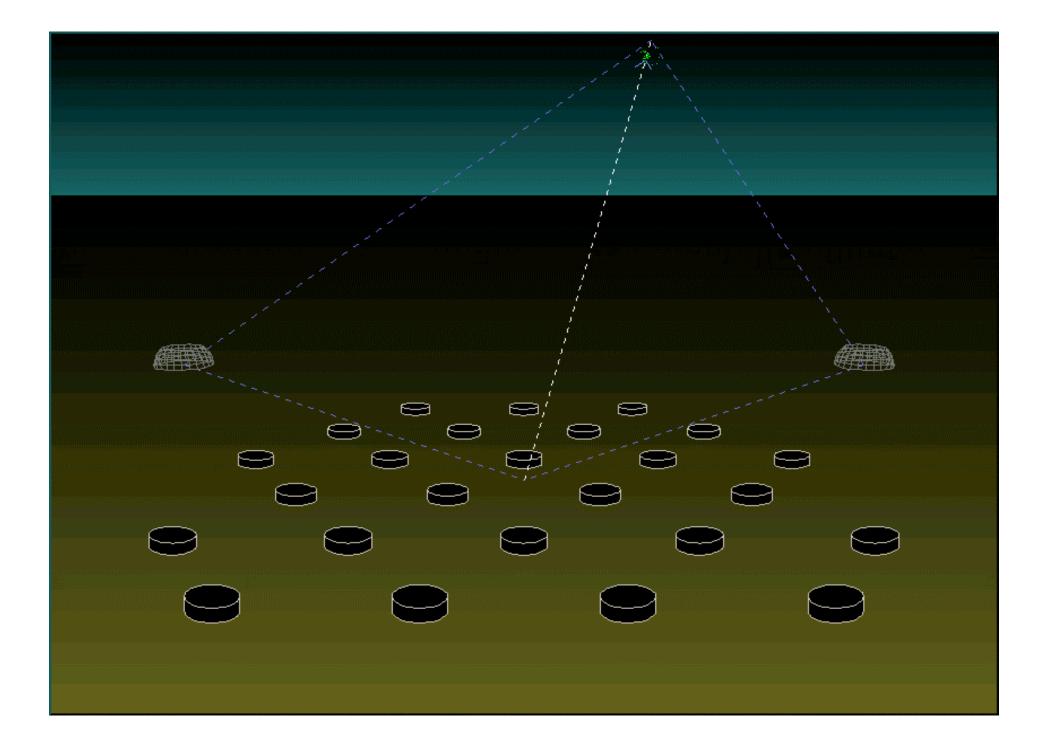
*Associate Countries

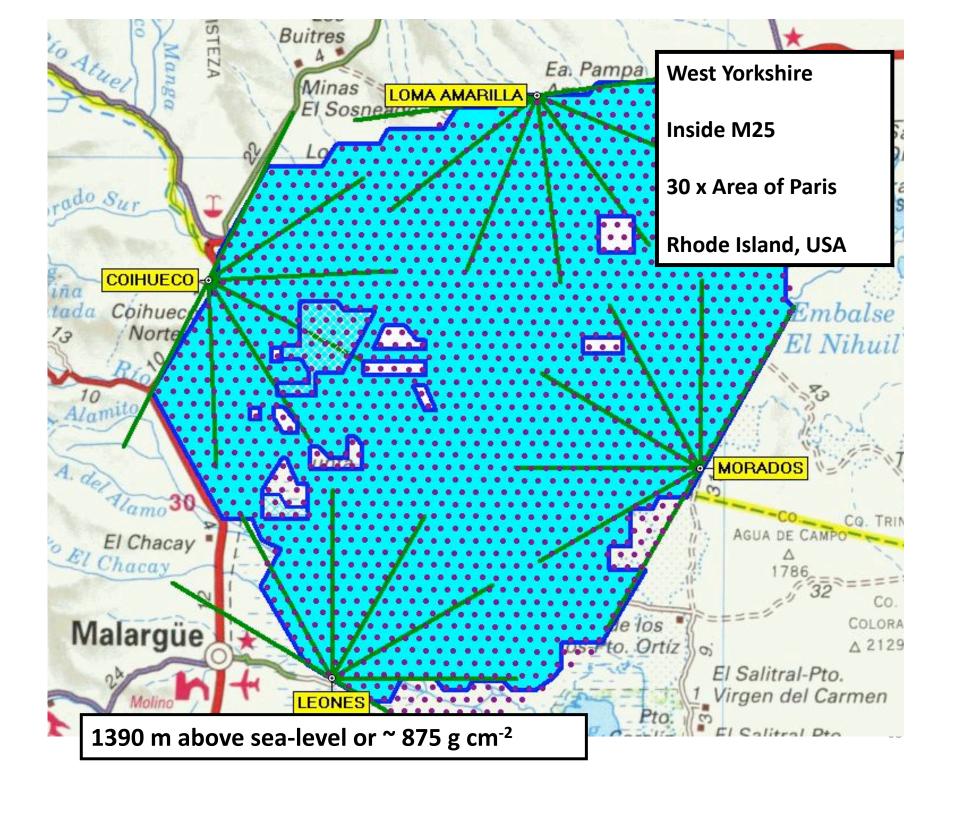
~ 400 PhD scientists from

~ 100 Institutions in 17

countries

Aim: To measure properties of UHECR with unprecedented statistics and precision – first discussions in August 1991 in Trinity College, Dublin – Jim Cronin and Alan Watson





Campus of Auger Observatory in Argentina

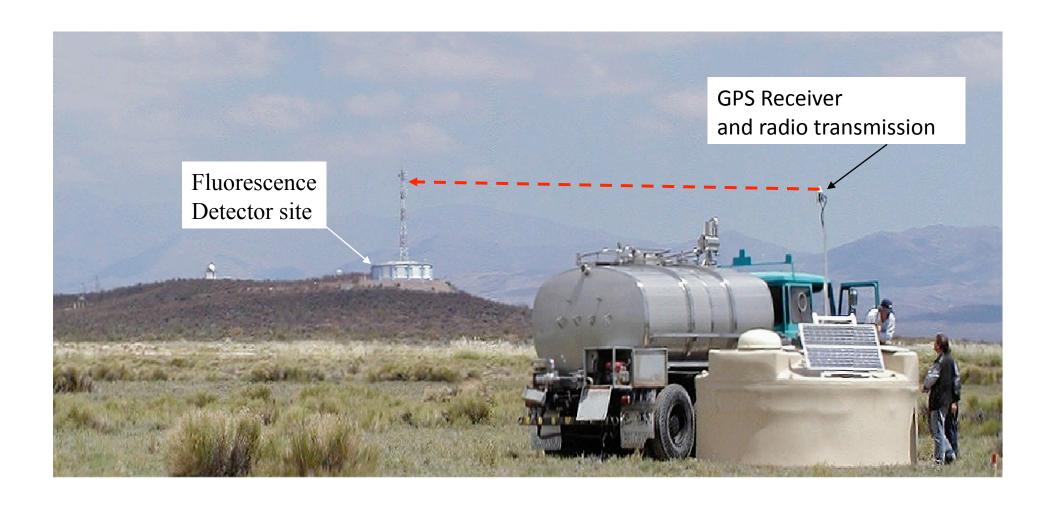




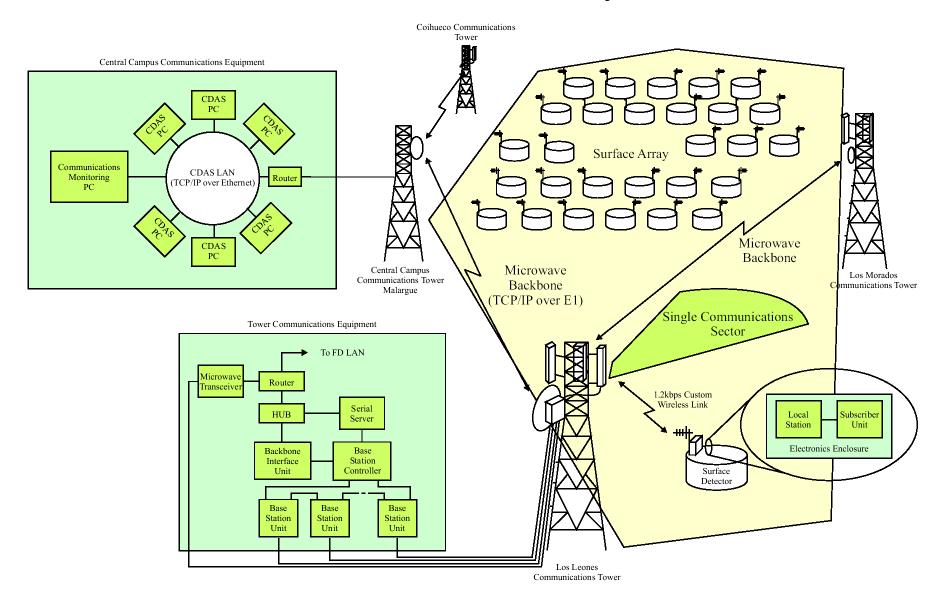




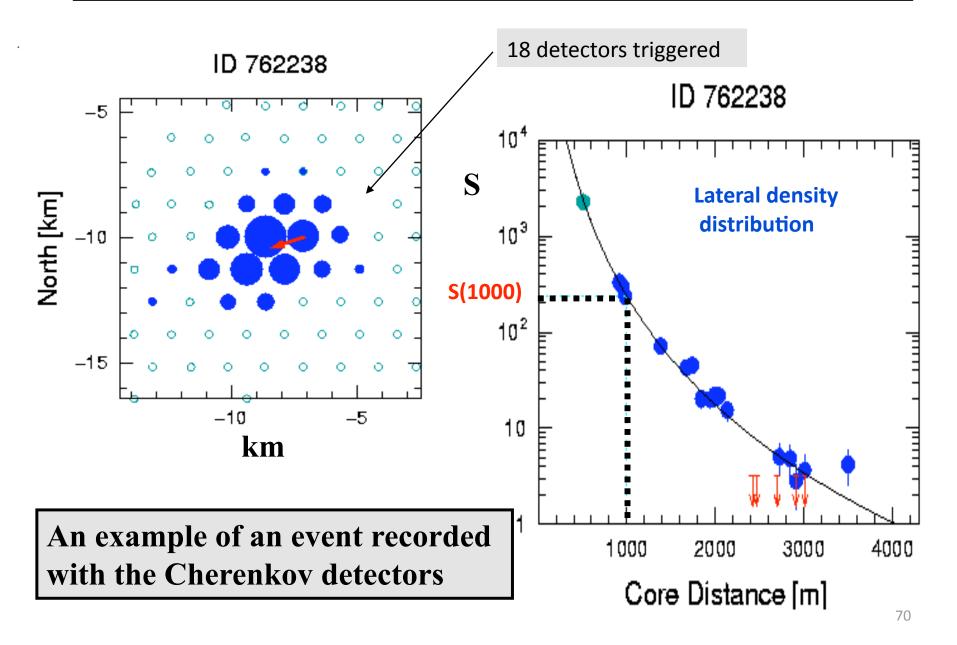
Last tank deployed: 13 June 2008

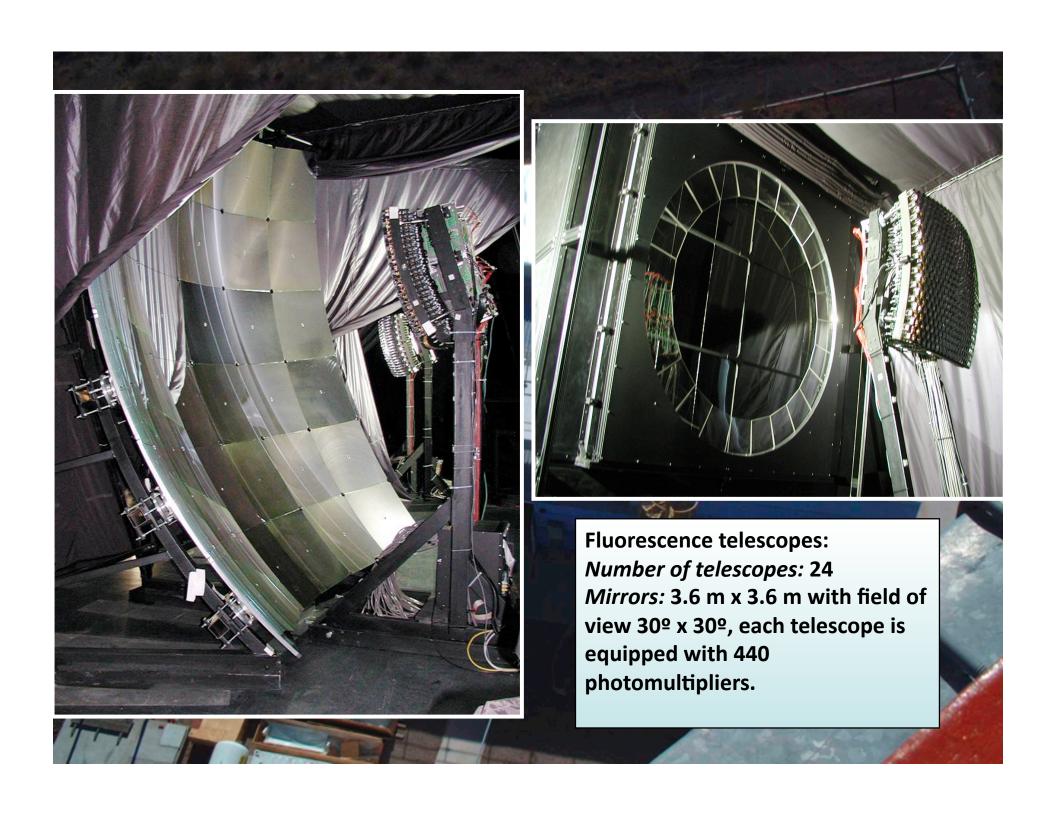


Telecommunication system



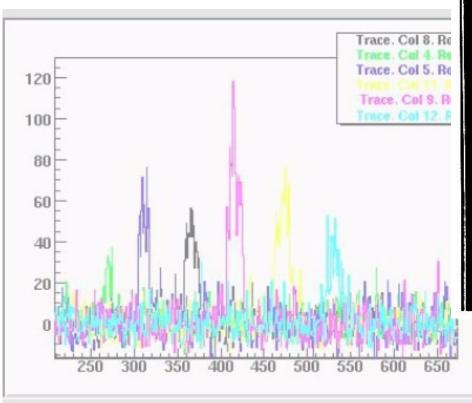
Zenith Angle ~ 48° Energy ~ 7 x 10¹⁹ eV

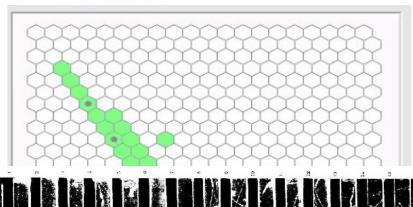


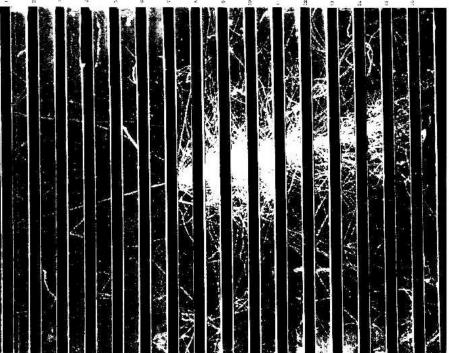


FD reconstruction

Signal and timing:-Direction and energy

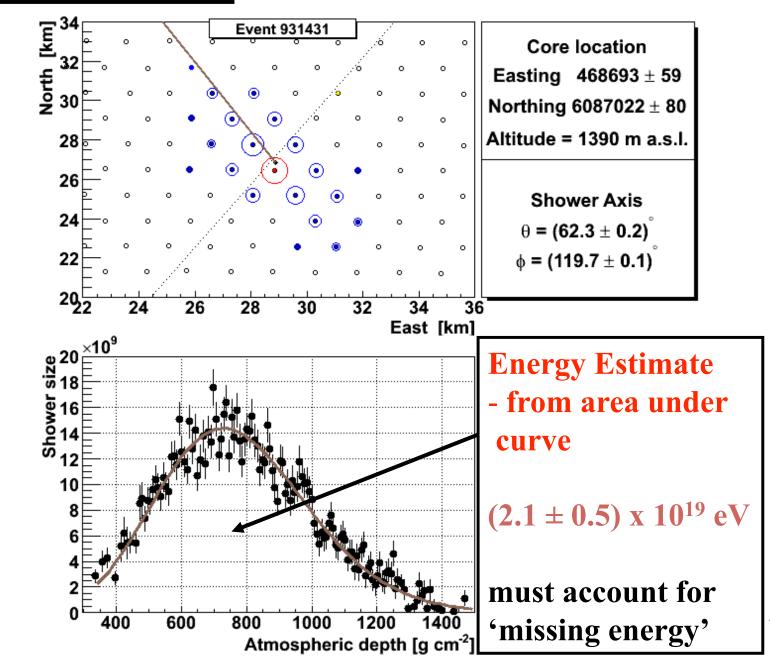


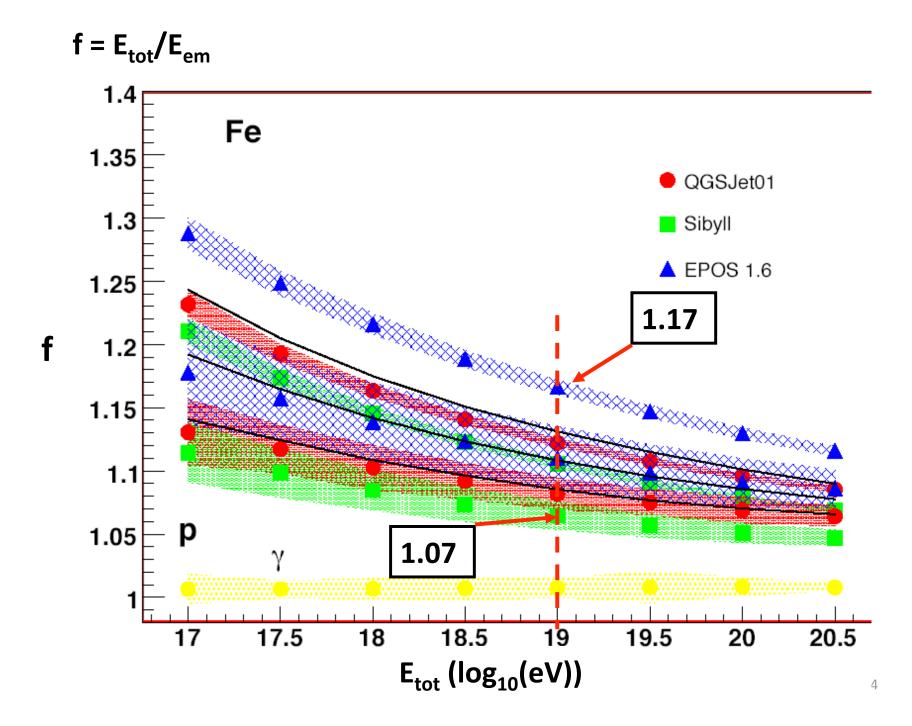




Pixel geometry shower-detector plane

A Hybrid Event





Results from Pierre Auger Observatory

Data-taking started on 1 January 2004 with

125 (of 1600) water-Cherenkov detectors

6 (of 24) fluorescence telescopes

more or less continuous operation since then

At end of 2009, 12,790 km² sr yr

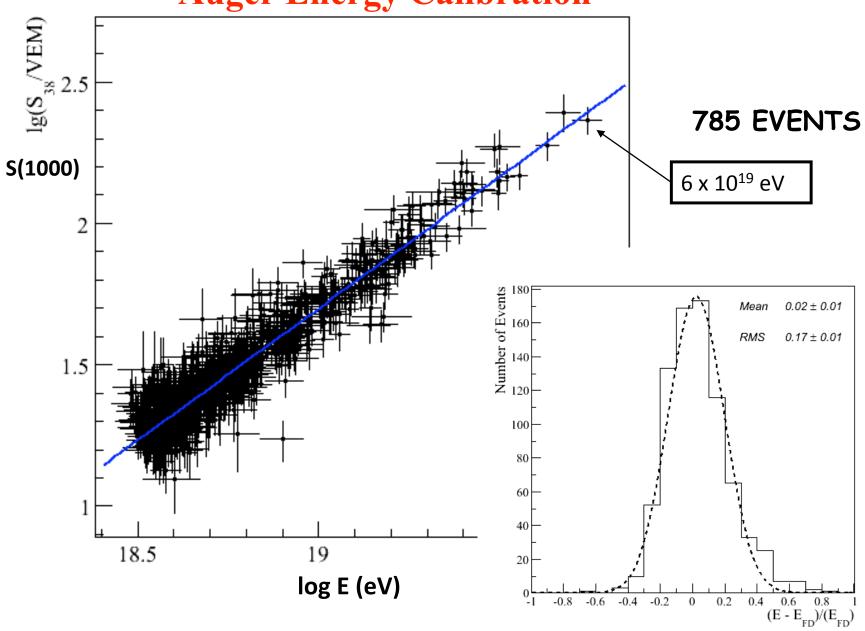
> 10¹⁹ eV: 4440 (HiRes stereo: 307

 $> 5 \times 10^{19} \text{ eV}$: 59 : 19

 $> 10^{20} \text{ eV}$: 3 : 1)

HiRes Aperture: x 4 at highest energies

Auger Energy Calibration

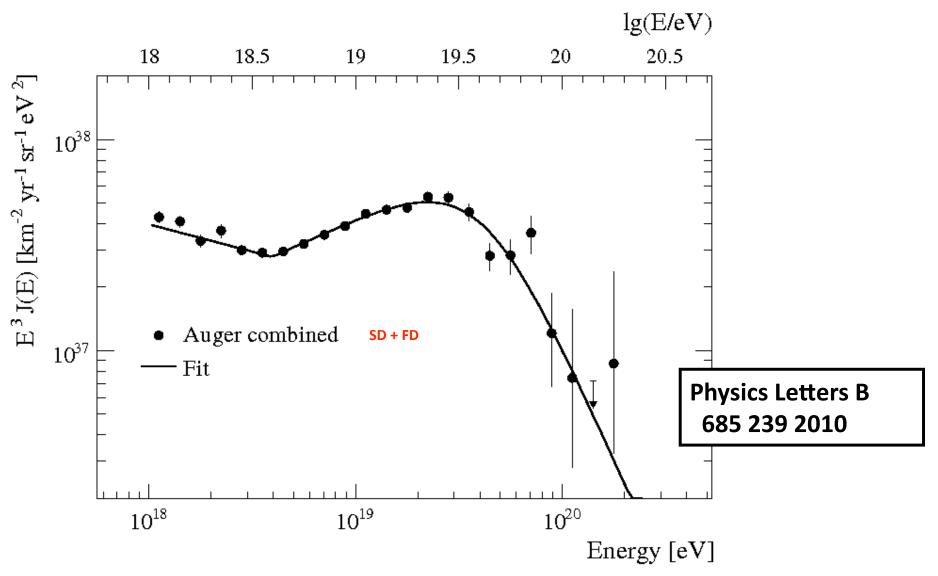


Summary of systematic uncertainties

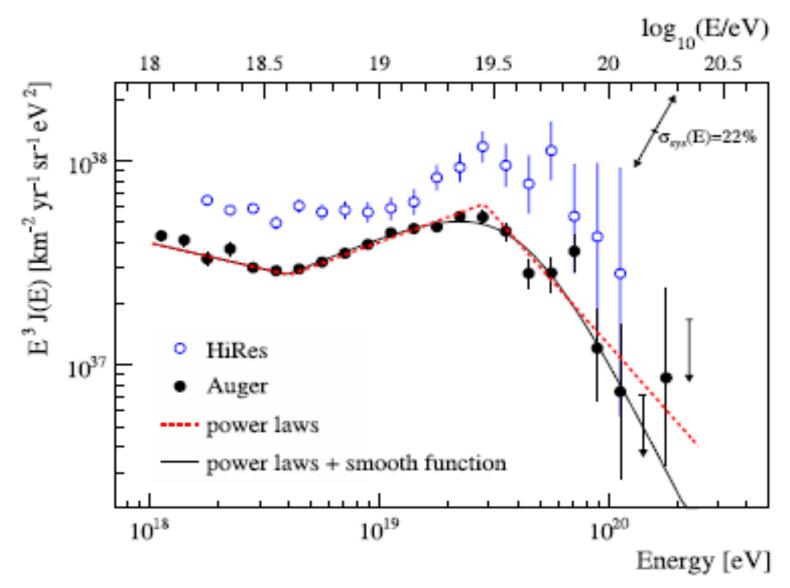
Source	Systematic uncertainty	
Fluorescence yield	14%	←
P,T and humidity	7%	
effects on yield		
Calibration	9.5%	←
Atmosphere	4%	
Reconstruction	10%	←
Invisible energy	4%	
TOTAL	22%	

Fluorescence Detector Uncertainties Dominate

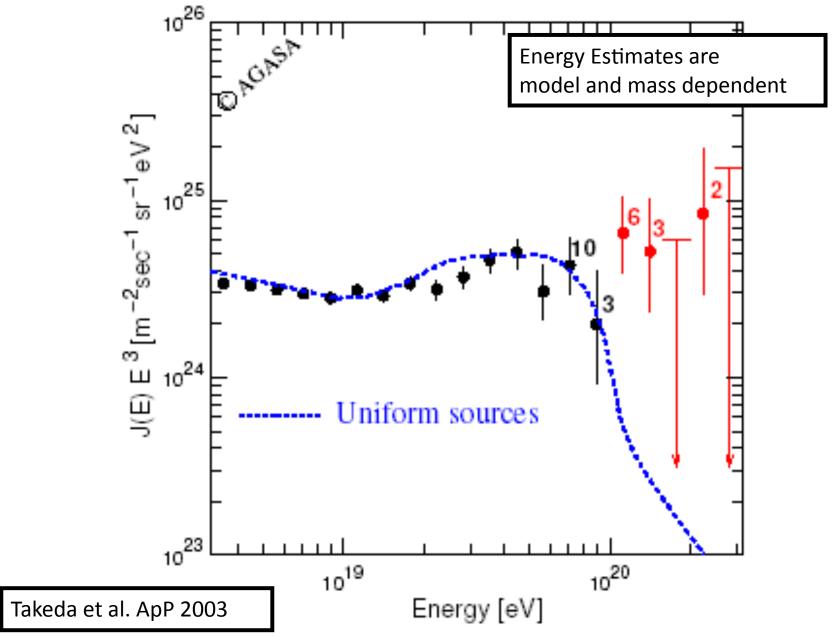
Energy Spectrum from Auger Observatory



Above 3 x 10¹⁸ eV, the exposure is energy independent: 1% corrections in overlap region



Auger and HiRes Spectra



For the few events above 10²⁰ eV

Auger (3) and HiRes stereo (1)

Integral flux is $(2.4 \pm 1.9/1.1) \times 10^{-4} \text{ km}^{-2} \text{ sr}^{-1} \text{yr}^{-1}$

11 AGASA events

 $(6.4 \pm 1.9) \times 10^{-3} \text{ km}^{-2} \text{ sr}^{-1} \text{ yr}^{-1}$

a factor of more than 25

Even a factor of x 2 increase in Auger energies would not be enough to explain difference

Consensus is that Auger and HiRes have got it right

But the steepening itself is **INSUFFICIENT** for us to claim that we have at last (predicted in 1966) seen the Greisen-Zatsepin-Kuz'min effect

It might simply be that the sources cannot raise particles to energies as high as 10^{20} eV – Nature could be teasing us!

But, if the steepening IS caused by the GZK-effect then we might expect to find that cosmic ray sources are relatively nearby

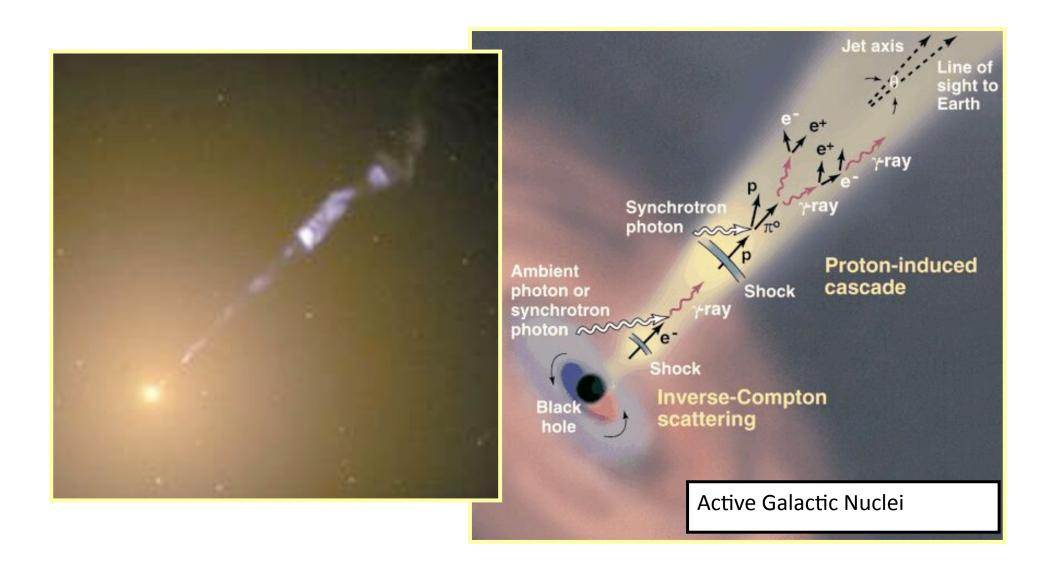
So, look to see if we can find likely accelerators lined up with the direction of the highest energy events

But what might the acceleration mechanism be?

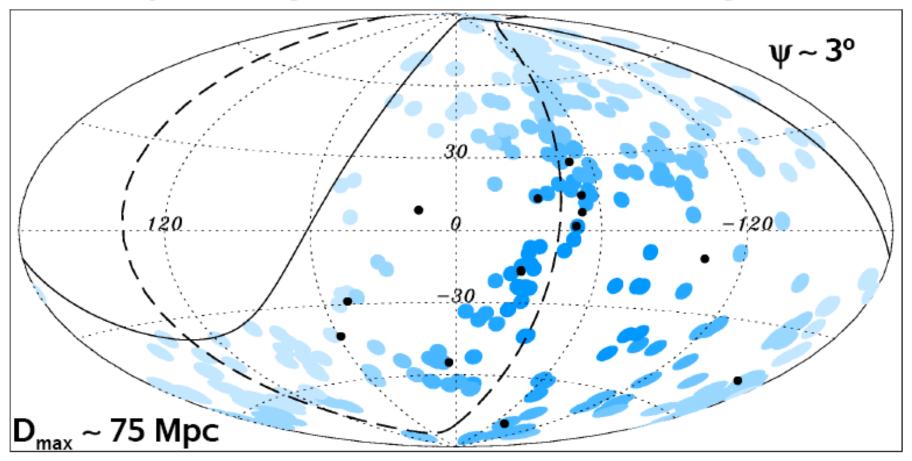
Searching for Anisotropies

Image of M87 with Hubble Space Telescope

Decided to use catalogue of galaxies like this and see if they lined up with the directions of our event



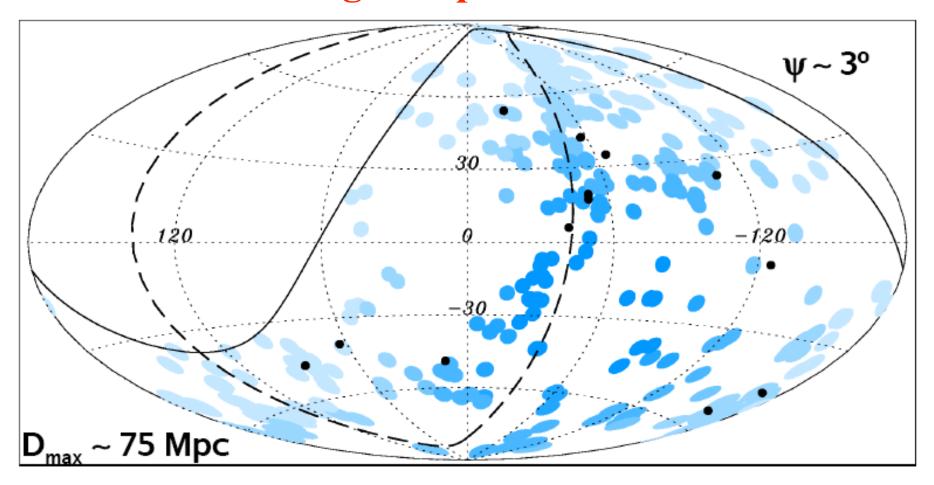
Exploratory scan: data until 27 May 2006



Largest significance for $E_{th} \sim 6 \times 10^{19} \; eV \; \psi \sim 3^{\circ} \; D_{max} \sim 75 \; Mpc$

12/15 events close to AGNs in Véron-Cetty & Véron Catalogue

Test Using Independent Data Set



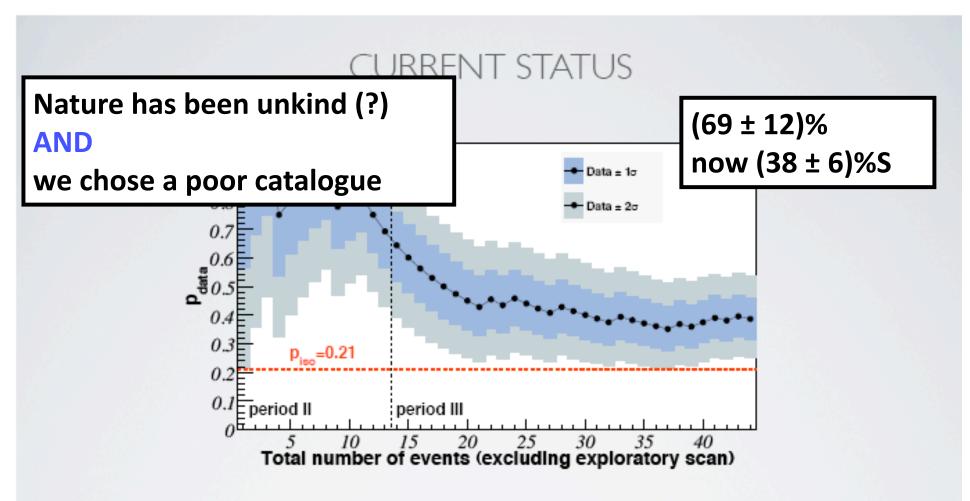
Data from 27 May 2006 until 31 August 2007 8/13 events lined up as before: chance 1/600

Using Veron-Cetty AGN catalogue

First scan gave ψ < 3.1°, z < 0.018 (75 Mpc) and E > 56 EeV

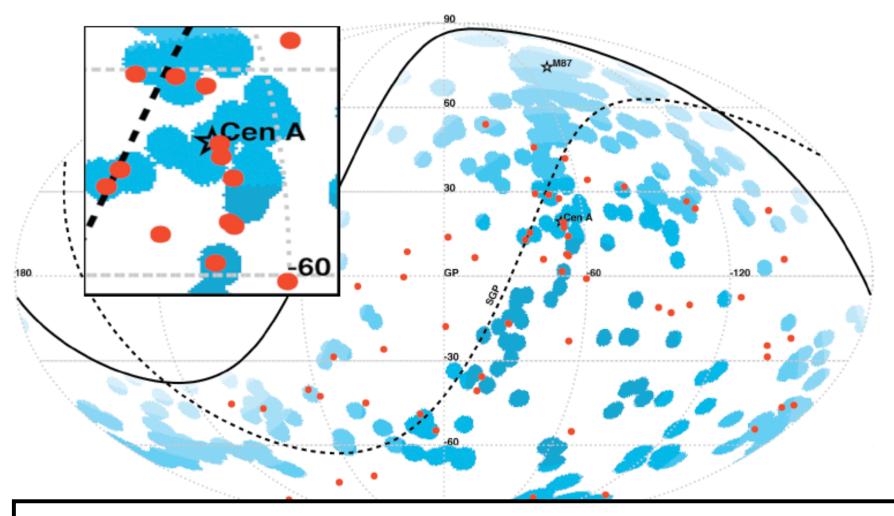
Period	total	AGN hits	Chance hits	Probability
1 Jan 04 - 26 May 2006	15	12	3.2	1 st Scan
27 May 06 – 31 August 2007	13 Each	8 exposure v	2.7 vas 4500 km²	1.7 x 10 ⁻³

6 of 8 'misses' are with 12° of galactic plane



p = 17/44 = 0.38 more than 2 s.d. from isotropy (expected from isotropy 9.2/44)

The degree of correlation has decreased, but still provides evidence for anisotropy of UHECRs @ E > 55 EeV at 99% C.L.

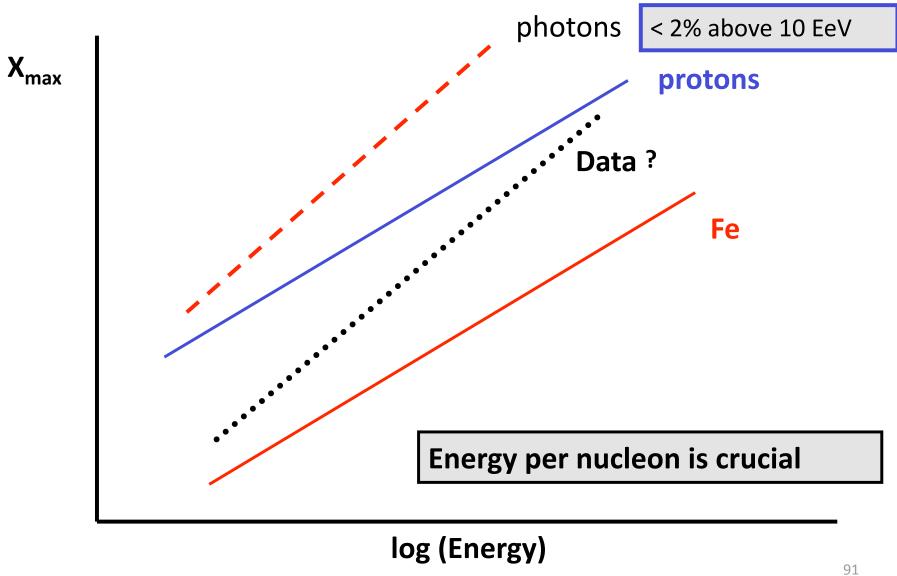


A clear message from the Pierre Auger Observatory is that we made it too small Rate of events that seem to be anisotropically distributed is only ~ 2 per month

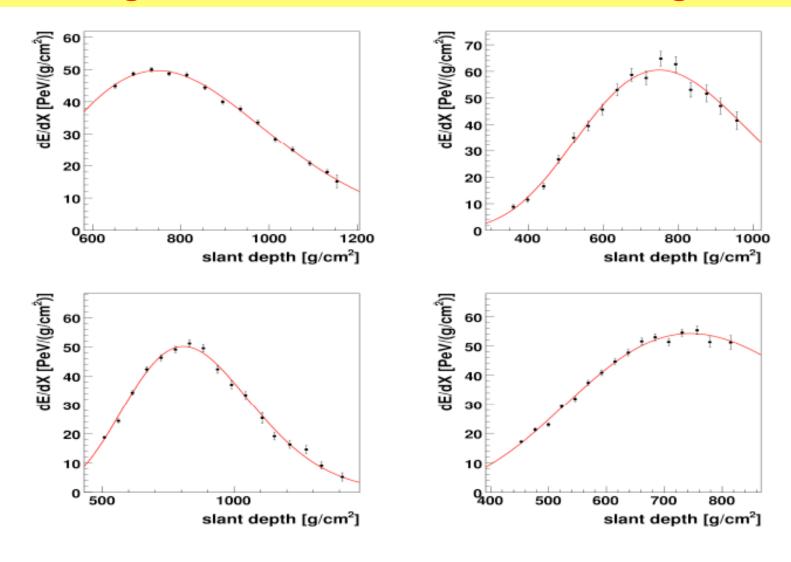
Indications on Mass Composition

- Anisotropy suggests a proton fraction of
 ≈ 40%
- Most unexpected result from Pierre Auger Observatory so far points in another direction
- Could be indicative of interesting new physics (??)

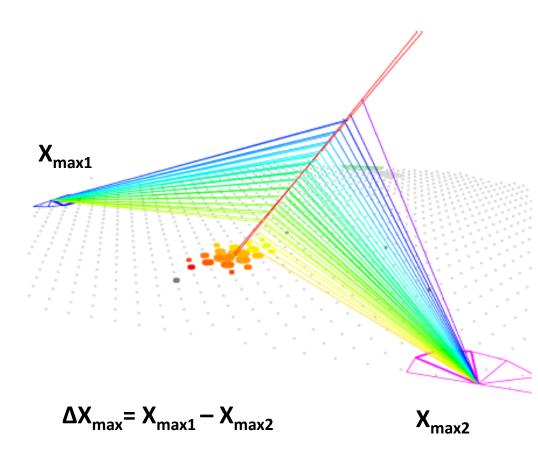
How we try to infer the variation of mass with energy



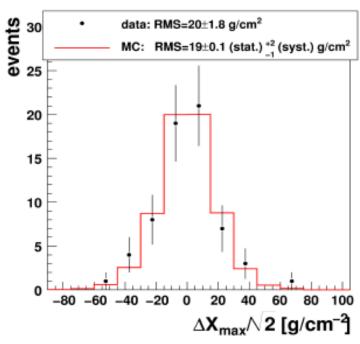
Some Longitudinal Profiles measured with Auger

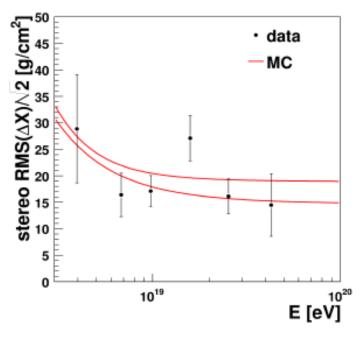


X_{max} Resolution

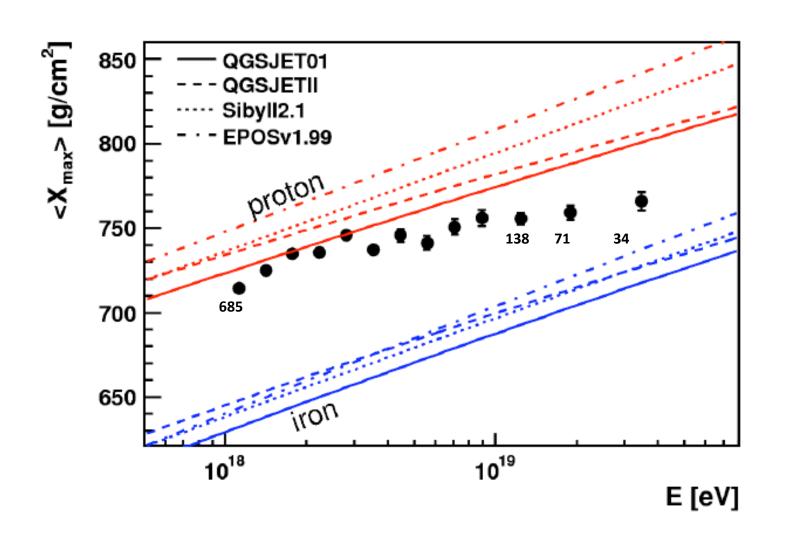


Check using Simulations

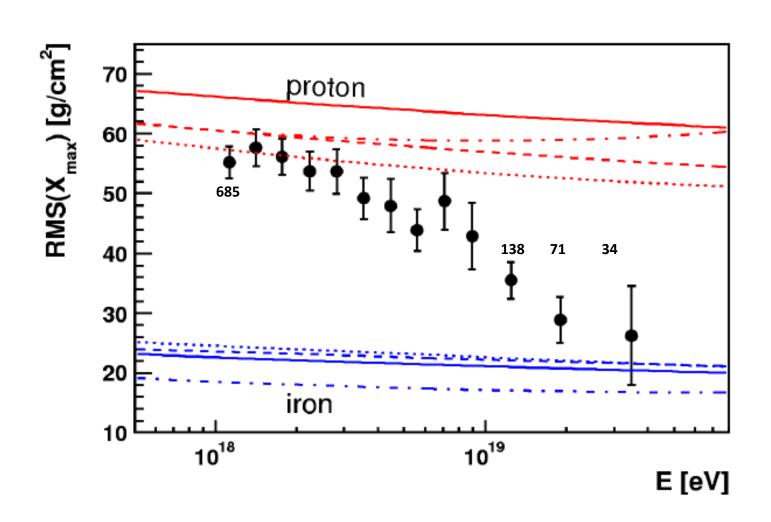




Mean X_{max} from 3754 events



RMS(X_{max)} for same events



Aloisio, Berezinsky and Gazisov Astroparticle Physics 34 620 2011

"Ultra High Energy Cosmic Rays: the disappointing model"

Assumes extra-galactic origin

Protons dominate 1 – 3 EeV

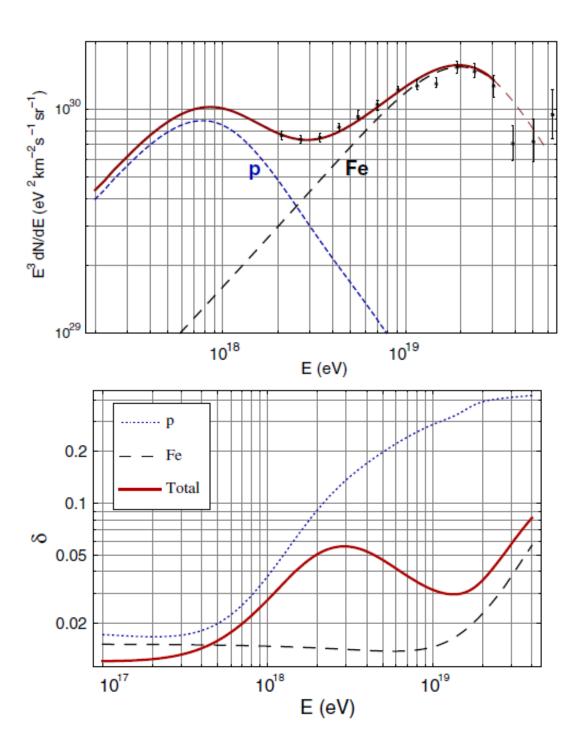
Low Maximum Energy of Acceleration: $E(max) = ZE_p$

No photo-pion production in Intergalactic space

GZK steepening does not exist

Absence of cosmogenic neutrinos

No anisotropy at high energies

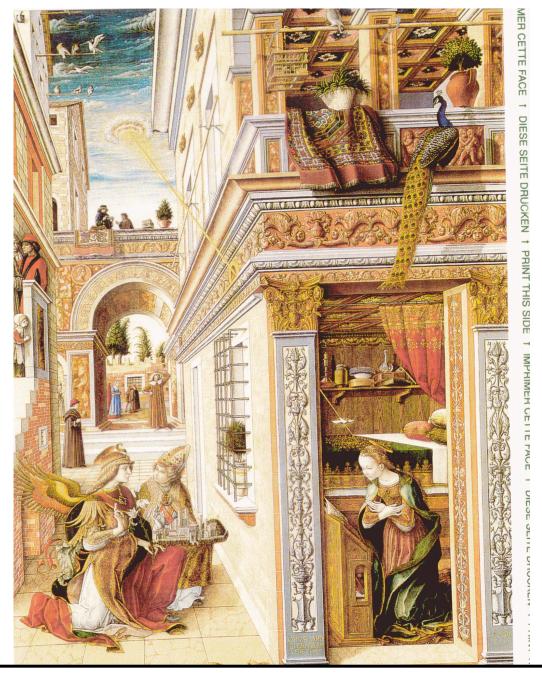


Calvez et al.

PRL 2010 105 09101

Consider a bursting source, GRB or rare types of SN explosions

Arrival Direction data cannot exclude this



Carlo Crivelli (1430 – 1490): 'The Annuciation with St Edimus'